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Complexity Classification of Product State Problems for Local Hamiltonians

arXiv: 2401.06725

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Definition

The k-LH problem is, given a k-local Hamiltonian, estimate its minimum eigenvalue f ground state energy.

This is analogous to the classical k-Max-SAT problem, where each clause acts on k variables.

We are often interested in the complexity of k-LH restricted to specific sets or families of Hamiltonians.

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S-LH classification

[Cubitt, Montanaro 2015], with [Bravyi, Hastings 2014], give a complete classification of 2-local S-LH as a function of S.

Given any set S of 2-qubit terms, [CM15] describes properties of the terms which determine whether S-LH is in P or NP-, StoqMA-, or QMA-complete.

What about product states?

What is the complexity of estimating minimum product state energies of various families of local Hamiltonians?

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- They're intermediate between classical states and general quantum states.
- For many natural sets of Hamiltonians, product states are rigorously near-optimal.

 $k\text{-}\mathsf{LH} o \mathsf{prodLH}$ $\mathcal{S}\text{-}\mathsf{LH} o \mathcal{S}\text{-}\mathsf{prodLH}$

 $k ext{-LH} o \mathbf{prodLH}$: given a $k ext{-local Hamiltonian, calculate the minimum energy over all product states: <math>\min_{|\psi\rangle} \langle \psi|H|\psi\rangle$ for $|\psi\rangle = |\psi_1\rangle|\psi_2\rangle\dots|\psi_n\rangle$.

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Can we classify the complexity of the product state problem for various families of Hamiltonians?



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Main Theorem (S-prodLH classification)

For any fixed set of 2-qubit Hamiltonian terms \mathcal{S} , if every matrix in \mathcal{S} is 1-local then \mathcal{S} -prodLH is in P, and otherwise \mathcal{S} -prodLH is NP-complete.

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Corollary

For any fixed set of 2-qubit Hamiltonian terms S, the problem S-LH is at least NP-hard if and only if S-prodLH is NP-complete.



For $\mathcal S$ any fixed set of 2-qubit Hamiltonian terms, if every matrix in $\mathcal S$ is 1-local then $\mathcal S$ -prodLH is in P, and otherwise $\mathcal S$ -prodLH is NP-complete.

Proof sketch



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- **To Do:** show if S contains a nontrivial 2-qubit term, then S-prodLH is NP-hard.

As an example, consider the 2-qubit term

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Let

$$|\phi_{\mathbf{v}}\rangle\langle\phi_{\mathbf{v}}|=\frac{1}{2}\left(I+v_1X+v_2Y+v_3Z\right)$$

Then the energy of the interaction between qubits u and v is

$$\operatorname{Tr}\left(H \mid \phi_{u} \rangle \langle \phi_{u} \mid \otimes \mid \phi_{v} \rangle \langle \phi_{v} \mid\right) = u_{1}v_{1} + u_{2}v_{2} + u_{3}v_{3} = u \cdot v$$

So for the example $S = \{X \otimes X + Y \otimes Y + Z \otimes Z\}$, the problem S-prodLH is equivalent to optimizing sums of inner products:

$$\sum_{uv\in E} w_{uv} \ u\cdot v$$

over unit vectors $u, v \in \mathbb{R}^3$.

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New goal: Given arbitrary 2-qubit *H*, does the optimum product state energy have a nice form like this? If not, can we force it to?

Write arbitrary 2-qubit H in Pauli basis:

$$H = \sum_{i,j=1}^{3} M_{ij}\sigma_{i} \otimes \sigma_{j} + \sum_{k=1}^{3} c_{k}\sigma_{k} \otimes I + w_{k}I \otimes \sigma_{k}.$$

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This is not as simple as $u \cdot v$, but we can design gadgets to simplify it.

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Trick 1: Symmetrize

It's nice when the objective function is symmetric, so acting on uv is the same as acting on vu.

Then we can work with un-directed graph problems.

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$$H_{sym} = H^{ab} + H^{ba} = H^{ab} + SWAP H^{ab} SWAP$$

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Use 4-qubit gadget with 2 ancilla

$$G^{uv} = H^{uv}_{sym} + H^{ab}_{sym} - H^{ua}_{sym} - H^{bv}_{sym}$$

$$u + v$$

$$- + -$$

$$b + a$$

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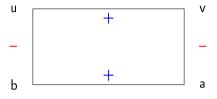
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Further analysis of gadget: $G^{uv} = H^{uv}_{sym} + H^{ab}_{sym} - H^{ua}_{sym} - H^{bv}_{sym}$



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$$\operatorname{Tr}(H |\phi_{\mu}\rangle\langle\phi_{\mu}| \otimes |\phi_{\nu}\rangle\langle\phi_{\nu}|) = u \cdot v.$$

Here, each edge/interaction H_{sym} also contributes

$$\operatorname{Tr}(H_{svm}^{uv}|\phi_u\rangle\langle\phi_u|\otimes|\phi_v\rangle\langle\phi_v|)\approx u^{\top}Mv.$$

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$$\operatorname{Tr}(H \mid \phi_{u} \rangle \langle \phi_{u} \mid \otimes \mid \phi_{v} \rangle \langle \phi_{v} \mid) = u \cdot v \approx 1 - \|u - v\|^{2}.$$

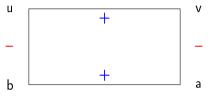
Here, each edge/interaction H_{sym} also contributes

$$\operatorname{Tr}(H_{\operatorname{sym}}^{uv}|\phi_u\rangle\langle\phi_u|\otimes|\phi_v\rangle\langle\phi_v|)\approx u^{\top}Mv\approx 1-\|Mu-Mv\|^2.$$

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Letting the ancilla a,b take optimal values, and summing the four contributions, we get

$$\|Mu - Mv\|$$

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We've used Hamiltonian gadgets to embed an objective function of the form

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Stretched linear Vector Max-Cut (MC_W^L)

For W a fixed diagonal matrix, and a graph G = (V, E), estimate

$$MC_W^L(G) = \frac{1}{2} \max_{\hat{u} \in S^{k-1}} \sum_{uv \in E} \|W\hat{u} - W\hat{v}\|$$

In words, assign unit vectors $\hat{v} \in \mathbb{R}^k$ to each vertex v in order to maximize the difference along each edge.

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Theorem

For any fixed non-negative nonzero $W = \operatorname{diag}(\alpha, \beta, \gamma)$ $\operatorname{MC}_W^{\mathsf{L}}$ is $\operatorname{NP-complete}$.

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Intuition: W defines an ellipsoid (if W=I, then its the unit sphere). Given some graph, the problem is to embed the vertices onto the ellipsoid's surface to maximize the sum of the edge lengths.

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Proof sketch:

1. Given a graph G, construct a new graph G' by replacing each edge with a 3-clique (triangle) gadget.

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- 1. Given a graph G, construct a new graph G' by replacing each edge with a 3-clique (triangle) gadget.
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This relates the MC_W^L value of G' to the 3-colorability of G. And 3-Coloring is NP-complete.



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 - 1. Construct Hamiltonian gadgets so the minimum product state energy has a nice form, like ||Wu Wv||.
 - 2. Show MC_W^L is NP-hard by a reduction from 3-Coloring.

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Quantum Max-Cut is equivalent to S-LH with $S = \{XX + YY + ZZ\}$.

Our classification theorem implies the following.

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3D-Vector-Max-Cut is NP-complete.

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Open problems:

- 1. Can we use complexity of product state problem to prove the *general* ground states of a class of Hamiltonians are *not* hard?
- 2. Classify S-prodLH with additional restrictions, e.g. only positive weights, spatial geometry?